| 1 | Pore-Pressure-Dependent Performance of Rocking |
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| 2 | Foundations |
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| 20 | Notion List | |
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| 21 | Α | Area of the footing |
| 22 | A _c | Minimum soil-footing contact area |
| 23 | α | Amplification |
| 24 | В | Width of the footing |
| 25 | D_w | Depth of water |
| 26 | D _r | Relative density |
| 27 | D_s | Depth of failure wedge |
| 28 | Δ | Horizontal distance between the centre of gravity and mid-foundation |
| 29 | Δu | Pore water pressure |
| 30 | f | Frequency of input motion |
| 31 | $f_{n,i}$ | Frequency at time instant <i>i</i> |
| 32 | h_{cg} | Distance from the footing base |
| 33 | K _r | Rotation stiffness |
| 34 | М | Mass |
| 35 | M _{c,foot} | The foundation theoretical ultimate moment capacity |
| 36 | L | Length of the footing |
| 37 | L _c | Critical length |

| 38 | PGA | Peak ground acceleration |
|----|------------------|--|
| 39 | r_s | Sphericity |
| 40 | r _r | Roundness |
| 41 | S | Degree of saturation |
| 42 | ξ | Damping ratio |
| 43 | ξ_i | Damping ratio at time instant <i>i</i> |
| 44 | $T_{n,i}$ | Period at time instant <i>i</i> |
| 45 | $	heta_f$ | Rotation of footing |
| 46 | W _{sat} | Saturated water content |
| 47 | ω_n | Natural frequency of oscillation |
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56 Abstract

57 This study explores the concept of rocking foundations to mitigate damage during seismic 58 events. By weakening the footing intentionally, the foundation acts as a "fuse" to prevent plastic 59 hinges from forming in columns. Shake table experiments were conducted on a lightweight 60 prototype deck mass-column-footing model founded on a fine, medium-dense sand, in two states 61 of nearly dry and saturated. Kinetic energy dissipation, hysteresis, and decay are examined for 62 various structure masses, for two nominal low and high motion frequencies. Findings suggest that 63 energy dissipation is higher in saturated sands — as compared to nearly dry and despite the absence 64 of liquefaction — due to fluctuating pore water pressure and a suction effect that evolves beneath 65 foundations' edge. Where the substrate retains its original porosity, flow of water and subsequent 66 damping becomes pivotal mechanisms to enable the rocking motion. In dry sands, energy 67 dissipation occurs mainly through rotation, and enhances with motion frequency (from testing 3 68 to 5 Hz). In saturated sands, energy dissipation occurs predominantly through plastic settlement, 69 and becomes less effective with motion frequency. Lighter structures experience greater rotational 70 movement, especially in the case of dry sands. That enhanced rotational movement is drove by 71 lower rotational stiffness. Overall, lighter structures facilitate the re-centring of the foundation 72 upon rocking motion.

73 Key words

74 Rocking; Rotation; Settlement; Saturation; Mass; Frequency; Damping

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78 INTRODUCTION

79 The recommended response for foundations, as per current geotechnical codes of practice (e.g., 80 EC8), remains to be elastic [1]. Conventional foundations are fixed base [2] and allow ductile 81 plastic hinging to form at the base of the column. However, an inelastic soil-footing response under 82 seismic shaking is inevitable and the consequent non-linear soil-structure interaction can become 83 beneficial in design. Seismic isolation through rocking motion is a relatively recent design 84 approach that provides seismic resilience by allowing subgrade soil to locally fail and enter the 85 non-linear plastic range of strains [3,4]. Foundation rapidly reaches its ultimate capacity, rotates 86 back and forth, and minimises the inertia forces that otherwise would transmit onto the structure.

In nature, the rocking effect resembles how tumbleweed responds to multilateral forces and moments. Tumbleweeds are detached from their roots and can rotate on their toes, lift, settle back down, and roll in windy weather. Fig. 1(a) shows the movement of a tumbleweed. Fig. 1(b) illustrates the rotations that cause the rocking foundation effect. The gravitational forces and breakaway gaps enable the system to restore the vertical orientation [5].

92 Seismic isolation

Rocking foundation as a concept was first introduced by Housner [6]. The concept incorporates weakening at the base through decreasing footing length or reinforcements. Consequently, the foundation theoretical ultimate moment capacity ($M_{c,foot}$) is decreased, transforming the foundation into a "fuse" to avoid plastic hinges forming in the column. The weakening naturally leads to a small factor of safety (2 to 8) for the vertical bearing pressure which then costs excessive foundation settlements [7]. Such movements necessitates incorporation of performance-based seismic design (PBSD), albeit PBSD in the rocking foundation context remains to be much further 100 studied [8]. Countermeasures include the placement of reaction blocks and concrete pads around 101 the foundation edges [9] piles and stone columns in softer soils [10]. The moment capacity $(M_{c,foot})$ is the primary factor that determines the energy dissipation. Gajan and Kutter [11] 102 103 formulated the moment capacity as a function of vertical load applied to the centre of gravity (P, 104 applied at a h_{cg} distance from the base), and the rocking re-centring distance (Δ , horizontal 105 distance between the centre of gravity and middle of soil-footing reduced contact area). The Δ is 106 associated with the contact area between the soil and the foundation, and how it evolves during 107 the motion. To this end,

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$$M_{c,foot} = P\Delta = \frac{PL}{2} \left(1 - \frac{A_c}{A} \right)$$
(1)

109 where A_c is the critical contact area needed to carry the vertical load of structure, thereby 110 dependent on factor of safety. In Fig. 1(b), L_c is the critical length, or the soil-footing contact length 111 over which the bearing pressure becomes equal to the ultimate bearing resistance as defined by 112 Terzaghi [12]. The L_c is the minimum length of footing-soil interface as foundation rocks. In Eq. 113 1, L is the length of footing along the direction of motion, and A is the footing area.



Fig. 1. Fuse mechanism: (a) tumbleweed plant detached from root and its survival from wind
effect through movement and minimised ductility demand (Adobe Stock #449307469, licence
AE02295300753CGB); (b) rocking motion as plastic hinge develops at soil-footing interface.

The predominant assumption is that foundations can rock when weakened, and that rocking takes place on homogeneous sands [13] that are fully saturated, dense, and well compacted [14]. Note that rocking is also a mechanism in clayey, cohesive soils, although it has been comparatively less studied. The PhD work of Hakhamaneshi [15] is a seminal contribution in this area. Their research showed that plastic settlements in clay are 20 to 40% less than typical settlements in sand under similar critical contact area ratios and rotation demands. They also linked the phenomena of breakaway gaps and uplift, as observed in rocking foundations, to those on sandy soils. Sharma
and Deng [16] referred to field snap-back trials and demonstrated a "better than on sand" recentring
capability of rocking foundations on clays (see also Sharma and Deng [17]).

127 In much of the previous literature, shallow foundations are also assumed to rest either at the 128 surface or with minimal embedment [23]. Sand is mostly considered dry [7, 11, 18, 19, 20] or 129 saturated [14, 20, 21, 22]. Groundwater level however fluctuates, leading to changes in the degree 130 of saturation of the subgrade soil. The impact of such variations has not received equal attention. 131 Burland and Burbidge [24] investigated the effect of the water table on foundation settlement and 132 constrained that association to a marginal 13% (in total settlement) for sands. Nevertheless, they 133 emphasized on the more significant overall effect on the ultimate bearing capacity. Under seismic 134 conditions, the water table influences the settlement, rotation, overturning moment, and energy 135 dissipation characteristics of rocking foundation systems [25]. The rocking foundation response 136 for unsaturated soils is not an objective in this paper, but this is an interesting, rapidly growing 137 field of study. Immediate expectations from unsaturated soils, when compared to their dry and 138 saturated state are an increase in soil stiffness and moment capacity, lower damping, and lower 139 foundation seismic settlement, and rotation at the surface [16, 23].

This paper presents the findings from experiments conducted on a model surface shallow foundation founded on a medium dense sand in two states of nearly dry and fully saturated. Resting on sand were four lightweight model columns and decks with a 0.04 scale factor. Mass was varied. A 1-g shake table was used to model a 0.2 g peak ground acceleration (PGA) earthquake at two, nominally low and high, 3 and 5 Hz frequencies to study the evolution of pore water pressure, energy (acceleration) dissipation, stiffness degradation, and displacements beneath foundation.

| 146 | Principles behind the fuse mechanism, design of the physical model, observations, and analysis |
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| 147 | are presented in the following sections. |
| 148 | EXPERIMENTAL PROCEDURES |
| 149 | Design principles |
| 150 | A mid-rise lightweight prototype deck mass-column-footing model was designed based on a |
| 151 | fixed-base 1.82-m diameter 11-m high column on Iranian standard F161 sand. Four design criteria |
| 152 | were followed: |
| 153 | (a) the system was a single-degree-of-freedom model with length scale factor of 0.04; |
| 154 | (b) the system was designed for a nominally elastic response of the column; |
| 155 | (c) a moment-to-shear ratio of $h_{cg}/B = 3$ was adopted, where h_{cg} is the height from the base |
| 156 | of the footing to the centre of gravity of the seismic mass and B is the length of the footing |
| 157 | in the direction of shaking. Theoretically, a $h_{cg}/B > 1$ moment-to-shear ratio encourages |
| 158 | rocking motion instead of translational movements; |
| 159 | (d) The critical contact area ratio, A/A_c , of the footing ranged from 1.81 to 2.69 regardless of |
| 160 | the soil saturation conditions. Here, A is the initial soil-footing contact area and A_c is the |
| 161 | minimum soil-footing contact area that supports the vertical load during rocking. The A_c |
| 162 | was determined in correspondence with method proposed by Deng and Kutter [26] and the |
| 163 | Brinch Hanson bearing capacity equations [27]. In Fig. 1(b), $L_c = A_c/B$. |
| 164 | Generally, in reduced-scale physical modelling, accurately reproducing stress fields within the |
| 165 | soil is challenging due to scale effects. These effects arise because the behaviour of soil under load |
| 166 | is pressure-dependent, which means that the smaller scale of the model may not perfectly mimic |

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the actual stress conditions in a full-scale scenario. The length scale factor of $\lambda = 0.04$ adopted

here means every 1 metre in the model represents 25 meters in the prototype. Despite the factor, the constrained depth of the soil layer inevitably leads to different interaction dynamics between the foundation and the soil compared to a full-scale model where soil depth might be more extensive. This includes potential differences in how the soil densifies or how stress is distributed through the soil column during seismic events.

173 Design principles led to a system with properties summarized in Table 1.

174 **Table 1** Model properties in scaled dimensions

| Property | Value |
|--------------------------------------|-----------------|
| Moment-to-shear ratio, h_{cg} /B | 3 |
| Critical length, L_c : m | 0.085 - 0.129 † |
| Critical contact area ratio, A/A_c | 1.81 - 2.69 † |
| D_w/B | 0.26 ‡ |

[†]Depending on mass of structure

176 ^{$\ddagger D_w$} is depth of water above pore water pressure transducers, *B* is footing width in the direction of shaking

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178 Model structure

179 A 3000 × 2000 mm ($L \times B$), indoor, one-dimensional shake table with 6 tonnes ultimate load 180 capacity and ±100 mm horizontal displacement range was used to generate uniaxial excitation (see 181 Fig. 2(a) and (b)). The table was designed and fabricated at Tabriz University. The foundation soil 182 was contained in a 650 × 1000 × 600 mm ($L \times B \times H$) rigid box composed of clear and stiff 183 PMMA panels and an exterior steel support frame. From bottom, the box was lined with a layer of 184 gravel which was in receipt of water through two tubes connected to drainage valves installed on the box sidewalls. A protective nonwoven geotextile layer covered the gravel and extended along the sides of the box. The rigid box was loaded onto the shake table using a crane (Fig. 2(c)) and subsequently bolted to the shake table base (Fig. 2(d)). Model footing was square in shape and 230 mm in width, resting on surface and above the subgrade sand. The footing was connected to a 230 × 230 × 82 mm rigid steel plate (i.e., load stub) via a 76.2-mm-diameter tubular column (Fig. 2(e)). The load stub provided seating space for four mass blocks (20, 25, 30 and 60 kg). The total mass of footing, column and load stud equalled 19.26 kg. The h_{cg} was set at 438 mm.



193 **Fig. 2**. Experimental set up: (a) shake table; (b) hydraulic actuator connected to control system;

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(c) installation of box on shake table; (d) and (e) deck mass-column-footing model.

195 Soil material

196 The foundation soil was a well-sorted ($C_u = 1.87$), fine (0.80 to 1.18 mm, and a $D_{50} = 0.27$ mm), and sub-angular ($r_s = 0.60$, $r_r = 0.38$) sand, similar to standard Ottawa C109 Sand in size 197 198 and shape. The sand is commercially available under F161 brand name and has a specific gravity of $G_s = 2.65$, maximum void ratio of $e_{max} = 0.815$ and minimum void ratio of $e_{min} = 0.534$. 199 200 Two sand conditions were examined: fully saturated and nearly dry. Sand was loaded into the box 201 in loose lifts and compacted with a hand-held vibratory compactor and a $580 \times 410 \times 50$ mm 202 plate, into 50 mm thick layers to achieve a target relative density of 75%. Compaction water 203 content varied between 2 to 5 wt.%. Based on the results of sand cone tests, conducted in various 204 stages of construction, the achieved mean relative density was 72.3%, corresponding to a void 205 ratio of 0.74, bulk unit weight of $\gamma = 15.71$ kN/m³ and a degree of saturation between 7.2 and 206 17.9%. The peak drained angle of friction for sand was 33° based on direct shear test results. For 207 saturated models, water was introduced to compacted soil in the box from two ports on sidewalls 208 of box, and through tubes that extend into the gravel layer. Valves were closed when water table 209 raised about 50 mm above soil surface. In saturated state, soil had a bulk unit weight of $\gamma = 19.12$ kN/m³ and a saturated water content of $w_{sat} = 27.8\%$. 210

211 Instrumentation

An array of 10 sensors, including five 0.01g in precision accelerometers (Acc), 2 linear variable differential transformers (LVDTs), three 50-kPa in capacity pore water pressure transducers and two cameras (for control) were utilised in experiments. Figure 3 shows a schematic front elevation view of the model and the instrumentation layout. One accelerometer was placed at the shake table 216 platen (Acc5), one on the mass block (Acc4), and three within the subgrade soil. Pore pressure 217 transducers were installed within the wedge of failure (also known as zone of shear and depth of 218 the rocking foundation) below the edges of the foundation and at a proximity. Rocking motion 219 takes place within this wedge, with depth of about $D_s = 0.5L_c.tan(45 + \emptyset'/2)$ — this is 220 approximately 20 mm in the present study. LVDTs were installed aligned with the direction of 221 seismic motion. Two cameras were used to capture videos of test sequence.

222 Testing diet

The test protocol included sinusoidal signals. Acceleration time histories are presented in Fig. 4(a) and (b). In total, 16 shake table experiments were performed: 8 experiments on nearly dry sand, constituting two motion frequencies of 3 and 5 Hz and four concentrated masses. Experiments were repeated for saturated sand.





Fig. 3. Schematic diagram of the physical model and the instrumentation layout



Fig. 4. Acceleration time histories for input signal: (a) f = 3 Hz, (b) f = 5 Hz.

231 ENERGY DISSIPATION

232 From Amplification Perspective

In general, greater energy dissipation in a substrate during vibration can be indicated by two readings. Firstly, evidence in support of reduced peak acceleration: If the peak accelerations are lower, one may conclude that the energy introduced by the vibration dissipates more within the medium. Secondly, evidence for increased damping characteristics: This can be observed if the vibrations die down more quickly, meaning the acceleration returns to zero faster.

238 Figure 5 illustrates a set of acceleration time history graphs, divided into two main categories 239 based on their location: beneath the foundation centre and at the top of the structure. Each set 240 includes four diagrams. Figs. 5(a), 5(b), 5(c) and 5(d) represent the acceleration measured by Acc2, 241 beneath the foundation centre. Figs. 5(e), 5(f), 5(g) and 5(h) illustrate the acceleration measured 242 by Acc4, at the top of the structure — that is near the centroid of the deck. Measurements were 243 constrained to two motion frequencies of 3 and 5 Hz. The case of saturated sand indicates lower 244 peak accelerations beneath and above the structure. The dissipation appears to be relatively more 245 rapid for systems subjected to higher frequency motion of 5 Hz.







The more marked energy dissipation in the case of saturated sand can be, partly, attributed to the relatively greater damping ratio.

252 Mechanisms

253 It will be shown, in Fig. 6, that motion generates an excess pore water pressure in saturated 254 sand. It will be also discussed that, motion-induced pore water pressure causes an array of vertical 255 displacements including formation of a breakaway gap. The latter enables the rotation, rocking, 256 thereby rapid fluctuations of pore water pressure across negative and positive values. As motion 257 matures, the pore water pressure subsides from its initial high due to gradual expulsion of water 258 from the soil network of pores. The substantial settlements in the case of saturated sand - see 259 Figs. 8(a)-(d), is consistent with effects of progressive water expulsion. The efflux of water is not 260 immediate despite the fundamentally drained nature of medium dense sand. In particular, in 261 regions close the centre (and below) of the foundation, the offset from drain points delays the 262 outflow of water. The flow of pore water to drain points creates internal friction within the sand, 263 leading to partial dissipation of energy. The measured lower amplifications (see Fig. 5) in saturated 264 sand is due to the latter damping effect. It is also evident that the ground motion amplification 265 modestly decreases with motion time, particularly at higher 5 Hz frequency. The decrease is, by-266 and-large, more pronounced in saturated sand. This progressive loss of amplification lends further 267 evidence to greater damping in saturated sands. Figs. 5(e) and (f) illustrate the registered 268 acceleration on top of the structure. As of the substrate, lower amplification was registered in structures on saturated sand. 269

270 From Pore Pressure Perspective

Figs. 6(a)-(b) show the variation of pore water pressure over time beneath the foundation. In Fig. 6(a), the rotational rocking motion is evident from the differences in pore pressures beneath the edge of the footing (T2) compared to points further away from the edge, yet still near the foundation (T3).

275 In principle, a registered positive pressure coincides with settlement, and a negative pressure 276 with uplift movement. During motion and for saturated sand, motion generates an immediate 277 excess pore water pressure, a drop in effective stress and a consequent buoyancy force with a 278 tendency to lift the foundation. The uplift causes a breakaway gap beneath the foundation, and 279 consequently development of negative pore water pressure. The negative pressures create a suction 280 effect that facilitate both movement of fluids and re-centring of the foundation. In Fig. 6(a), 281 negative pressure appears to be more marked under the edge of the footing (T2) and retains 282 amplitude for the entire motion duration (i.e., 6 seconds). As such, for any structure mass, it 283 appears that the breakaway gap remains open to enable water influx and efflux. In contrast, the 284 negative pressure decreases with motion time in surrounding areas (see T3). Upon termination of 285 motion, pore pressures tremble within a positive range of values before returning to hydrostatic 286 levels. This indicates post-excitation expulsion of water, densification of soil (i.e., consolidation), 287 and ground settlement. Note the presence of positive pore pressures (on termination of the 288 vibration) leads to plastic settlements after dynamic motion. This will be further discussed in 289 following sections.

Fig. 6(b) illustrates the variation of pore water pressure with time beneath the edge of the footing and for two concentrated masses of 20 and 60 kg. Beneath the edge of the footing and under the lightest structure (20 kg), for the select motion time range of 4.2 to 6.2 s, Figs. 6(c) and 6(e) show the variation of motion acceleration and pore water pressure respectively. Figure 6(d)

294 and 6(f) show similar diagrams for the edge of the footing under the heaviest structure (60 kg). 295 Beneath the lightest structure (20 kg) and edge of footing, pore water pressure appears to tremble 296 more frequently. This is a manifestation of relatively easier re-centring and can be explained from 297 two perspectives. Firstly, the relatively more porous subgrade soil beneath the lighter structure 298 provides sufficient space for pore water to mobilize. The short wavelength cycles in Fig. 6(e) may 299 be attributed to that greater porosity. Secondly, the relatively more permeable substrate beneath 300 the lighter structure facilitates formation of the breakaway gap, which then accommodates the 301 mobilized water and causes the tremors seen in Fig. 6(e). The gap is an enabler for the rocking 302 motion. The amplification beyond the input $\alpha = 0.2$ g on top of the structure in Fig. 5(g) and 5(h) 303 is consistent with this enhanced rotation. In Fig. 7, the range of footing rotation angles increase 304 with decreasing structure mass. In other words, lighter structures benefit from a breakaway gap to 305 experience greater levels of footing rotation. Regardless of structure mass, negative pore water 306 pressures retain amplitude for the entire motion duration.



Fig. 6. Saturated sand profiles (a) pore water pressure evolution beneath the footing's edge (T2)
and proximity (T3); (b) pore water pressure evolution beneath the edge of footing for varied
surcharge; (c) and (e): acceleration and pore pressure evolution beneath the footing – light
structure; (d) and (f): acceleration and pore pressure evolution beneath the footing – heavy
structure.

314 ROCKING RESPONSE

315 Hysteresis Response

- 316 Fig. 7 exhibits normalised overturning foundation moment versus rotation in the direction of
- 317 input motion (θ_f) for nearly dry and saturated conditions. De-centralised and modestly wider
- 318 hysteresis loops for saturated sands indicate enhanced levels of damping and energy dissipation.
- 319 This is consistent with acceleration time histories observed in Fig. 5.



Fig. 7. Normalised moment-rotation hysteresis loops

322 Irrespective of motion frequency and mass, saturated sands generally exhibit lower peak 323 moment and peak rotation (also see Antonellis et al. [28] and Turner et al. [25]). The response 324 seems to be influenced by factors such as the initial packing quality, applied vertical load, and 325 embedment depth. In this study, the subgrade sand used is relatively more porous (D_r of around 326 72%), supports relatively lighter masses, and underlies a surface footing without any embedment. 327 To this end, although the nearly dry sand develops some apparent cohesion (through matric 328 suction) to arrest rotation, the cohesion and subsequent restrictions to motion are strictly limited. 329 In Fig. 7, the hysteresis loops for nearly dry sand close with increasing structure mass, indicating 330 a reduction in energy dissipation. The motion results in a greater degree of densification of 331 subgrade soil beneath the heavier structure, a reduction pore space size and, consequently, an 332 increase in matric suction when conditions favour its development.

333 Total Kinetic Energy

321

For the mass-pier-foundation system subjected to dynamic excitation, the total kinetic energy constitutes the kinetic rotation energy and the kinetic translational energy. The kinetic rotation energy is equal to the contained area in each hysteresis loop. The kinetic translational energy is a factor of foundation settlement.

In Fig. 7, the nearly dry sand exhibits a much larger hysteresis area. In other words, energy dissipation in dry sands is principally through rotation and gains momentum with increasing motion frequency from 3 to 5 Hz. In Fig. 8, saturated sand exhibits a larger cumulative plastic settlement. In other words, energy dissipation in saturated sands is principally through settlement,

particularly under heavier masses, and loses momentum with increasing motion frequency from 3to 5 Hz.

344 **Displacements**

Fig. 8 illustrates the vertical displacements of the foundation. As one would expect, nearly dry sand develops relatively greater values of elastic settlement, whereas saturated sand develops relatively greater values of plastic (and cumulative) settlement. The latter is consistent with observations made in Fig. 7 and in the context of total kinetic energy, where the predominant mode of energy dissipation in saturated sands was established to be through plastic settlement.

The input motion frequency (within the limits of testing 3 and 5 Hz) is inversely related to both elastic and plastic settlement. However, this relationship is more pronounced in subgrade soils under sufficiently high surcharge loads.

The uplift movement of foundations in the form of ground heave is limited to lower 3 Hz loading frequency and early stages of cyclic loading; that early uplift is particularly notable in dry sands. This observation is consistent with earlier discussions, where large rotational movements of the foundation under the testing 3 Hz frequency motion were associated with uplift movements in dry sands.

In Fig. 8(f), a conceptual model of nearly dry sands subjected to shaking is presented. Capillary action causes water to move upward through the soil pores due to surface tension. During seismic excitation, sand particles compress and expand cyclically, leading to water being drawn upward through capillary action. This results in an increase in pore water pressure near the foundation base, a decrease in effective stress, and a reduction in the sand's ability to resist uplift forces. The combination of higher pore water pressure and lower effective stress facilitates the uplift



movement of the foundation in the early stages of motion, leading to rotation and rockingmovements.

Fig. 8. Vertical displacements: (a) subgrade soil beneath light structure subjected to 3 Hz motion,
(b) subgrade soil beneath heavy structure subjected to 3 Hz motion, (c) subgrade soil beneath
light structure subjected to 5 Hz motion, (d) subgrade soil beneath heavy structure subjected to 3
Hz motion, (e) concept of elastic and plastic settlements, (f) conceptual model of capillary action
and heave in nearly dry sand

372 Possible Limitations

373 As noted, in the provision of target relative density, a small amount of water was mixed with 374 dry sand ahead of compaction. In this, the dry sand in this study is in effect a nearly dry sand, 375 despite it being fair to assume, that large proportions of that small added water would have 376 evaporated ahead of the applied seismic motion. The existence of small amounts of water in dry 377 sand merits some elaborations on possible implications. The cumulative foundation rotations and 378 settlements reported here for the nearly dry sand were likely to be slightly larger, had the sand been 379 completely dry. Note that the rotational energy dissipation at the surface is through two coinciding 380 kinetic rotational and kinetic translational mechanisms. Thereby it is hard to say what would have 381 been the effect on the rotational energy dissipation had the sand been completely dry. Cautiously put, rotational energy dissipation at the surface is likely to be modestly lower, have the sand been 382 383 completely dry. The smaller rotations, in the presence of matric suction, albeit little, are due to 384 enhanced stiffness and strength.

385 Rotational Stiffness and Damping

386 In the context of soil dynamics, rotational stiffness (K_r) refers to the resistance provided by a 387 soil-structure system to rotation under the influence of applied moments or forces. It characterises the soil's ability to deform and rotate in response to external loads. Mathematically, K_r represents the rate at which the moment is changing concerning rotation.

390 The damping ratio (ξ) characterizes the level of damping or energy dissipation in a soil-391 structure system subjected to dynamic loads or cyclic loading. It describes how rapidly oscillations 392 in the system decay over time. A higher damping ratio indicates more significant dissipation and 393 quicker decay of vibrations.

394 Methods followed in Sharma and Deng [29] are adopted here. The rotational stiffness can be 395 represented by the slope of the line that passes through the opposite ends of the hysteresis loop on 396 the $M - \theta$ diagram. This is demonstrated in Equation 2.

$$397 K_{sec} = \frac{M_{(\theta_{max})}}{\theta_{max}} (2)$$

However, the maximum moment (M_{max}) does not necessarily coincide with the maximum footing rotation (θ_{max}) . Therefore, the ratio of the maximum moment in each cycle to the maximum recorded rotation in that cycle may provide a more accurate representation of rotational stiffness (see Equation 3). During cyclic loading, the contact surface between the footing and the subsoil may alter at any moment due to the footing's rocking motion, resulting in either an increase or a decrease in the contact area. Any change in the contact surface leads to a corresponding adjustment in the system's rotational stiffness. To this end, the secant stiffness is,

$$405 K_{sec} = \frac{M_{max}}{\theta_{max}} (3)$$

406 In Figs. 9(a)-(b), the stiffness is normalised by the initial stiffness, which is the slope of the 407 linear portion of the $M - \theta$ curve. 408 The damping ratio is calculated via a simple MATLAB code, and by dividing the area enclosed 409 by each $M - \theta$ hysteresis loop by four times π times the area of a triangle. This triangle has a 410 height of M_{max} and a base of θ_{max} . The formula is expressed as,

411
$$\xi = \frac{\Delta E_{cycle}}{4\pi\Delta E_{el}} \tag{4}$$

412 where ΔE_{cycle} is the area enclosed by a specific hysteresis loop and ΔE_{el} is the area of a triangle 413 defined by the coordinates (0,0), (0, θ_{max}), and (M_{max} , θ_{max}).

414 Figure 9(a) and (b) demonstrate the evolution of normalised rotation stiffness with loading cycles for the lightest and the heaviest structures, for the f = 5 Hz motion. Immediate inspection 415 416 shows that the normalised rotational stiffness decreases with loading cycles and increases with 417 mass of structure. The degradation is non-linear. Turner et al. [25] reported similar non-linearity 418 and attributed that to soil rounding beneath the foundation. The direct relationship between 419 normalised stiffness and mass marks a restriction on the rocking motion for heavy structures. The 420 lower foundation rotation angles (θ_f) in Fig. 7 for heavier structures correspond well with the 421 observations here. It also evident that the dependency of rotational stiffness on mass becomes more 422 marked in saturated sands.

Figures 9(c) to 9(f) demonstrate the variation of damping with loading cycles for the lightest and heaviest structures subjected to the f = 3 Hz and f = 5 Hz motions.

The damping ratio generally decreases with loading cycles. In other words, cyclic loading appears to compromise the soil structure. Observations here are consistent with findings of Khosravi et al. [5], Turner et al. [25] and Allmond [30] for sands at various saturation states. To explain the concurrent degradation of damping and rotational stiffness, two mechanisms can be 429 identified: (i) Due to cyclic loading, particles rearrange, and the packing becomes denser. As 430 particles draw closer together, the coordination number increases [31], leading to a rise in inter-431 particle contact points. This change facilitates a more uniform distribution of skeletal stresses, 432 resulting in a decrease in energy absorption per cycle. (ii) In saturated sands, excess pore pressures 433 lead to softening and a partial loss of the soil's capacity to dissipate energy through damping.

It is interesting to see the consistency of equivalent damping ratio for saturated sand with those of clay (0.08-0.3) in Sharma and Deng [16], up to 0.4 in Sharma and Deng [17], and viscous damping of 0.10-0.35 in Sharma and Deng [29]. When comparing the data, reader should note that, the horizontal axis in Fig. 9 reads the number of input cycles and that may differ from seemingly similar diagrams in the literature. In the context of pore water, viscous damping is relatively higher in saturated sands and is powered by flow of pore water through the pore network.





Fig. 9. Normalised rotational stiffness and damping versus loading cycles



443 To better identify the nonlinearity and rocking performance of testing systems, it is vital to measure and capture the variation with time of fundamental period $(T_{n,i})$ and damping ratio (ξ_i) , 444 445 and the associated system's acceleration response amplitudes. Time here refers to post-vibration 446 conditions, through which the acceleration decays. Of seminal works that address the damping of 447 shallow foundations subjected to rocking oscillations are Tomassetti et al. [32], Spanos et al. [33], 448 Adamidis et al. [34], Wiebe et al. [35], Anastasopoulos et al. [36], and Makris and Konstantinidis 449 [37]. Most recently, de Silva et al. [38] correlated the nonlinear variation of the foundation stiffness 450 and damping ratio with the increasing amplitude of the foundation translation and rocking motions (θ_{fmax}) and proposed sets of empirical equations. To this end, for two concentrated masses of 20 451 452 and 60 Kg and higher motion frequency of 5 Hz, the procedures proposed in Arabpanahan et al. [39], and Kashani et al. [40-41] were followed to determine the frequency at time instant ($f_{n,i}$ = 453 $1/t_{i+1} - t_i$). The t_i is the time instant i and the $t_{i+1} - t_i$ is the period, $T_{n,i}$. The procedure was 454 455 also utilised to determine the damping ratio at time instant i – that is ξ_i . The procedure is based 456 on the free decay vibration, when the input motion stops, and system freely vibrates to return to 457 static conditions and behaviour transitions from nonlinear to linear. Free decay motion can be seen 458 in the acceleration time history in Fig. 10(a). The motion amplitude subsides during the free decay 459 due to damping. The procedure constituted five steps. The free decay response was trimmed at 460 both ends through removing the static components. The acceleration response's low-magnitude, 461 high-frequency components were eliminated using a zero-phase fourth-order low-pass 462 Butterworth filter, which has a cut-off frequency of 10 Hz. Zero-phase filtering ensures that there is no time shift in the filtered signal. The amplitude at time instant t_i (for the i^{th} vibration) was 463 464 then determined $(A_{c,i})$. This led to the determination of the frequency and period at time instant t_i $(f_{n,i} \text{ and } T_{n,i})$. An exponential curve was then fitted to $A_{c,i}$ and $A_{c,i+1}$ — corresponding to time 465

466 instants t_i and t_{i+1} , to determine the damping ratio, ξ_i . The latter procedure is demonstrated in 467 Fig. 10(b) and Equation 5. In Figure. 10(b), to ascertain the instantaneous damping ratio (ξ_i), an 468 exponential curve is fitted to X_{ci} and X_{ci+1} at time instances t_i and t_{i+1} , respectively.

469
$$\log_e \frac{x_A}{x_B} = \frac{2\pi\xi_i}{\sqrt{1-\xi_i^2}}$$
 (5)



470

471 Fig. 10. (a) free decay motion on Acceleration response time history, (b) damping ratio (ξ_i) at time 472 instant t_i

Points A and B in Equation 5 stand for two consecutive peaks on an underdamped system's response curve. In the time instance of t_A and t_B , respectively, x_A and x_B denote the amplitude corresponding to points A and B, respectively.

476

Figure 11 demonstrates the free decay response for the lightest (M = 20 Kg) and heaviest (M = 60 kg) structures. The frequency at time instant t_i ($f_{n,i}$) increased with time and as the amplitude dropped progressively. The rise in frequency can be attributed to the decrease in rotation,

- 480 re-establishment of soil-structure contact and stiffness. Mass is directly correlated with the $f_{n,i}$,
- 481 indicating the more marked nonlinearity of soil behaviour under heavier structure.





Fig. 11. The motion decay for input 5 Hz sinusoidal motion, (a)-(b) period at varied time instances,
(c)-(d) frequency at varied time instances, (e)-(f) filtered and trimmed acceleration response, (g)(h) damping ratio at varied time instances.

For both dry and saturated sands, an increase in decay time (and hence decreasing acceleration amplitude) led to decreasing period (T_n) and reciprocal increase in frequency $(f_{n,i})$ at time instant t_i — see Fig. 11(a)-(d). This is consistent with earlier observations of Irani et al. [41-42] and indicative of foundation re-centring as the contact surface between soil and foundation reestablishes. The change in frequency with time is relatively faster in the case of saturated sands, indicating the impacts and contribution of evolving pore water pressures. Saturated sands experience relatively lower damping ratios at any given post-motion time.

493

494 CONCLUSIONS

495 Rocking foundation incorporates weakening of footing and compromising the foundation 496 moment capacity, so to transform the foundation into a 'fuse' to avoid plastic hinges forming in 497 the column. This study examines the rocking foundation on a fine, medium-dense sand beneath a 498 range of lightweight masses and explores the inter-playing of mass, motion frequency, and degree 499 of saturation. A single-degree-of-freedom, mid-rise lightweight prototype deck mass-column-500 footing model is built on sand and subjected to sinusoidal shaking.

501 Findings here are constrained to two testing input motions of 3 and 5 Hz frequency, and the 502 F161 well-sorted, fine, sub-angular sand (similar to Ottawa C109 in shape and size) compacted to 503 a relative density of just above 70% in two fully saturated and nearly dry states.

505 In the case of saturated sand,

| 506 | 1. | Lower peak accelerations and amplification beneath and above the structure, alongside |
|-----|---------|---|
| 507 | | decentralised and wide $M - \theta$ loops, suggest greater energy dissipation. |
| 508 | 2. | The principal mechanism of energy dissipation is plastic settlement, especially under |
| 509 | | heavier loads. |
| 510 | 3. | Negative pore water pressures, peaking beneath the foundation's edge, maintain |
| 511 | | amplitude throughout the motion, indicative of a breakaway gap with sustained porosity. |
| 512 | 4. | Rapid changes in pore pressures beneath the foundation's edge indicates the breakaway |
| 513 | | gap's role in facilitating rocking motions. |
| 514 | 5. | Water efflux towards dissipating excess pore water pressure, results in significant plastic |
| 515 | | settlements once the motion ceases. |
| 516 | 6. | Foundation recentring appears easier under relatively lighter structures, evidenced by |
| 517 | | more frequent changes in pore pressures beneath the foundation edge, a wider range of |
| 518 | | footing rotation angles, and larger amplifications above the lighter structure. |
| 519 | In | the case of nearly dry sand, |
| 520 | | 1. The $M - \theta$ loops close with increasing surcharge, indicating decreased energy |
| 521 | | dissipation as surcharge increases. |
| 522 | | 2. Energy dissipation primarily occurs through rotation. |
| 523 | | 3. Notable foundations uplift at the early stages of motion is consistent with large |
| 524 | | rotational movements. |
| 525 | For all | testing sands, |

| 526 | 1. | The normalised rotational stiffness and damping degrade nonlinearly across loading |
|-----|----|---|
| 527 | | cycles, possibly due to particle rearrangement and more uniform distribution of skeletal |
| 528 | | stresses. |
| 529 | 2. | During the free decay phase, both the period (T_n) and amplitude decrease, leading to |
| 530 | | progressive increase in frequency at time instant t_i ($f_{n,i}$). These changes are more |
| 531 | | pronounced in saturated sands and suggest foundation recentring as the contact surface |

- 532 between soil and foundation re-establishes.
- 533

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669 AUTHOR STATEMENT

Esmatkhah Irani, A. came up with the concepts, designed the methods, fabricated the test
setup, conducted experiments and analysis, and wrote the first draft. Hajialilue-Bonab, M came
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Assadi-Langroudi, A. conducted supplementary analysis, wrote the second draft, edited the
artworks, also the final copy. Maleki Tabrizi contributed to visualization, instrumentation, editing,
and data management.

676

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